

THE CMAF WINDOW

Possible sizeable energy production
from 500/1000 eV deuterons

Main topics covered

- I The physics and maths used in alpha disintegration constants and d/d fusion reaction rates. The Yukawa potential.
- II Determination of the coupling constant of the Yukawa potential (alpha disintegration)
- III Determination of d/d reaction rates
- IV Comparison with experimental data (Huke and SPAWAR).
- V The coupling resonance and the CMAF window.

I - Physics and Maths

The Gamow penetration factor γ

Plays a major role

- In alpha disintegration constants
- In d/d fusion reactions rates

$$\gamma = \frac{2\sqrt{2m}}{\hbar} \int_{R_1}^{R_2} \sqrt{(B(r) - E)} dr \quad (m = \text{reduced mass})$$

I - Physics and Maths

Expression of the barrier

- The barrier equation is:

$$B(r) = +ke^2 \left(\frac{1}{r} \right) + \frac{l(l+1)\hbar^2}{r^2} - k' Cg^2 \frac{e^{-r/\rho}}{r}$$

Coulomb Centrifugal Yukawa

- With

*$k = 1$ and $k' = 4$ in the d / d case
 $k = 2Z'$ and $k' = 4A'$ in the α case*

- Calculation of the tunnelling probability $P(E_d)$ using a spreadsheet

I - Physics and Maths

Expression of the barrier

- The boson carrying the Yukawa interaction might be a neutral and virtual electron/positron pair, with mass $2m_e$ (electron mass)

(A. Meulenburg suggestion August 2008)

- Its range would be:

$$\rho = \hat{\lambda} = \frac{\hbar}{2m_e c} = 193 \text{ fm}$$

- The coupling constant C can be calculated from known experimental values of the alpha disintegration constants.

II - Alpha disintegration case

- Following relations allow the determination of C (Yukawa coupling constant) by fitting calculated alpha disintegration constants λ with measured ones:

$$\lambda = \nu P$$

$$\nu = \frac{1}{2R_1} \sqrt{\frac{E_\alpha}{2m_\alpha}}$$

$$P = e^{-\gamma}$$

$$C = 3.79 * 10^{-6} g^2 = 7.53 * 10^{-3} e^2$$

III - d/d reaction rates case

$$r(\text{cm}^{-3}\text{s}^{-1}) = \sigma(\text{cm}^2)\varphi(\text{cm}^{-2}\text{s}^{-1})N_0(\text{cm}^{-3})$$

σ cross section $\sigma = \sigma_{\text{geom.}} P(E_d)$ probability of tunnelling

N_0 deuteron concentration in palladium

φ incident deuteron flux ($\text{cm}^{-2}\text{s}^{-1}$)

● The energy produced is:

$$W = r(1 + 3 + 0.83 + 2.44)1.6 * 10^{-13} \text{ (W cm}^{-3}\text{)}$$

III - d/d reaction rates case

- The d/d reaction rates were calculated (with $P(E_d)$ given by the model) for an incident deuteron flux corresponding to 1 W with energy E_d (varying from 2 to 100,000 eV) on a 1 cm² target, containing 1 mmole of d (d/Pd = 0.7)
- Calculations were run in 2 cases : no screening and screening + action of the Yukawa potential (with strength C from the alpha case).
- The influence of the Yukawa was found negligible at a few eV (huge impact of the screening) and of the same order of magnitude as the electron screening at a few keV.

III - d/d reaction rates case

- With these hypothesis, the number of incident deuterons is:

$$n = \frac{1}{100E_d} \sqrt{\frac{m_d}{2E_d}} \text{ (cm}^{-3}\text{)} \quad (E_d \text{ in eV)}$$

and their flux:

$$\varphi = \frac{1}{E_d} \text{ (cm}^{-2}\text{s}^{-1}\text{)}$$

The palladium target thickness is $126 \mu\text{m}$,
containing 1 mmole d ($6 \cdot 10^{20}$ d)

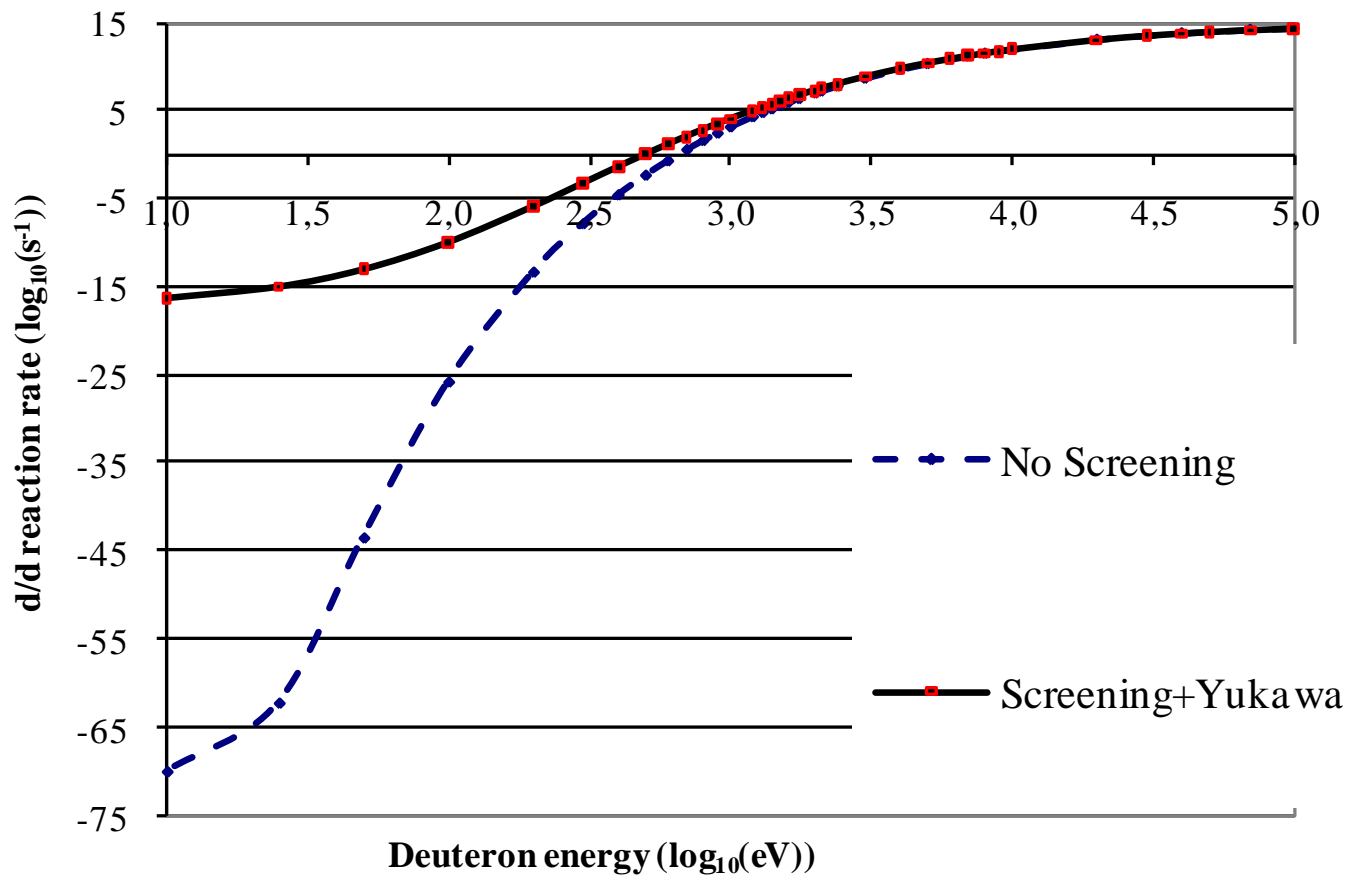
corresponding to $N_0 = 4.74 \cdot 10^{22} \text{ (cm}^{-3}\text{)}$

IV - Comparison with experimental data

Calculated d/d reaction rates

Figure 1

d/d reaction rates for 1 W in

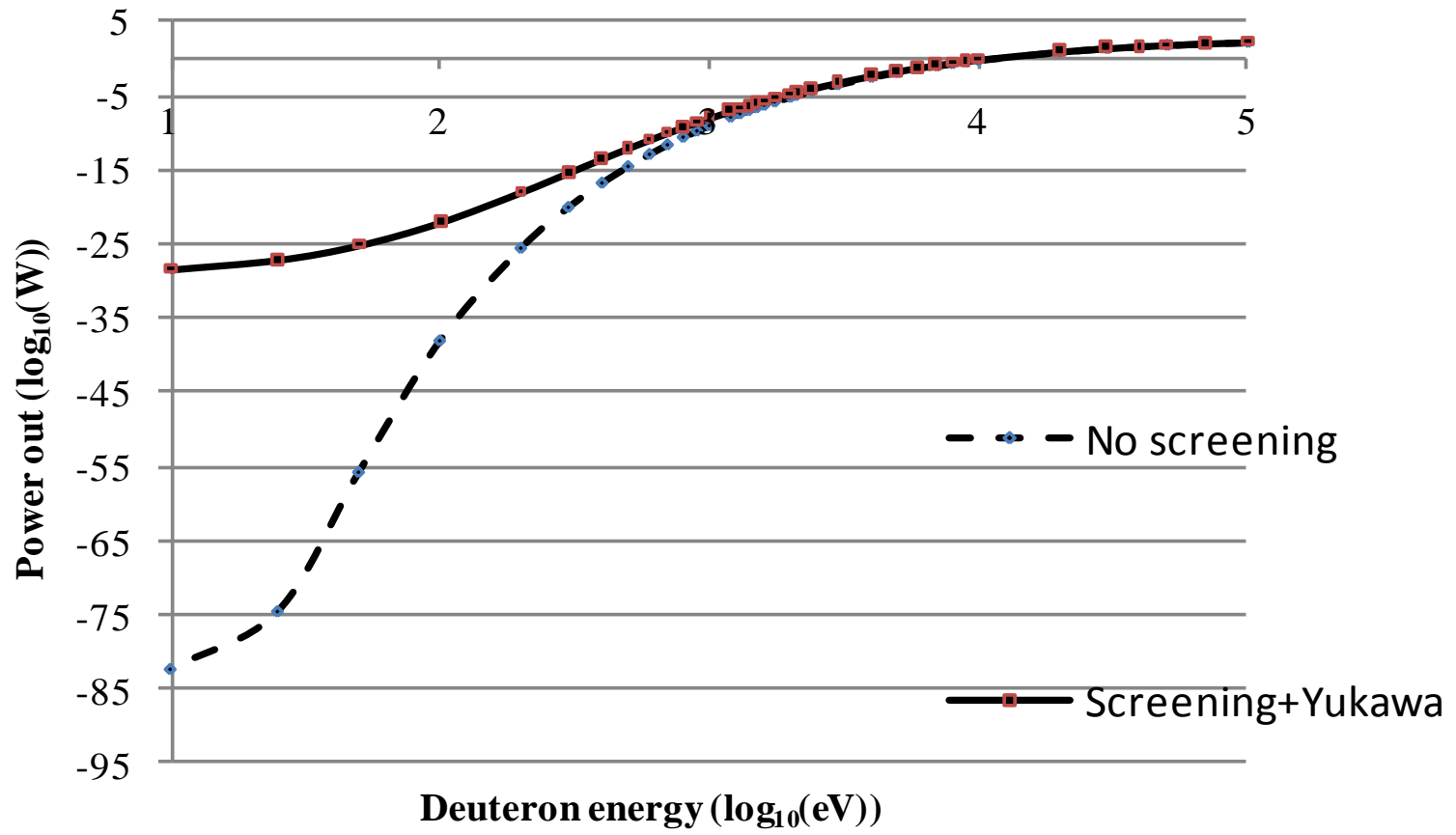


IV Comparison with experimental data

Calculated Power out

Figure 2

Power out for 1 W in



IV - Comparison with experimental data

- Hucke results (ref.1) show experimental reaction rates r_{ex} higher than calculated ones r_{cal} , with screening and Yukawa (ref.3)

$$r_{ex} = F_c r_{cal}, \text{ with } F_c = 1.5 \text{ to } 2$$

- SPAWAR results (ref.2) show experimental reaction rates r_{ex} very much higher than calculated ones r_{cal} , with screening and Yukawa (the latter negligible)

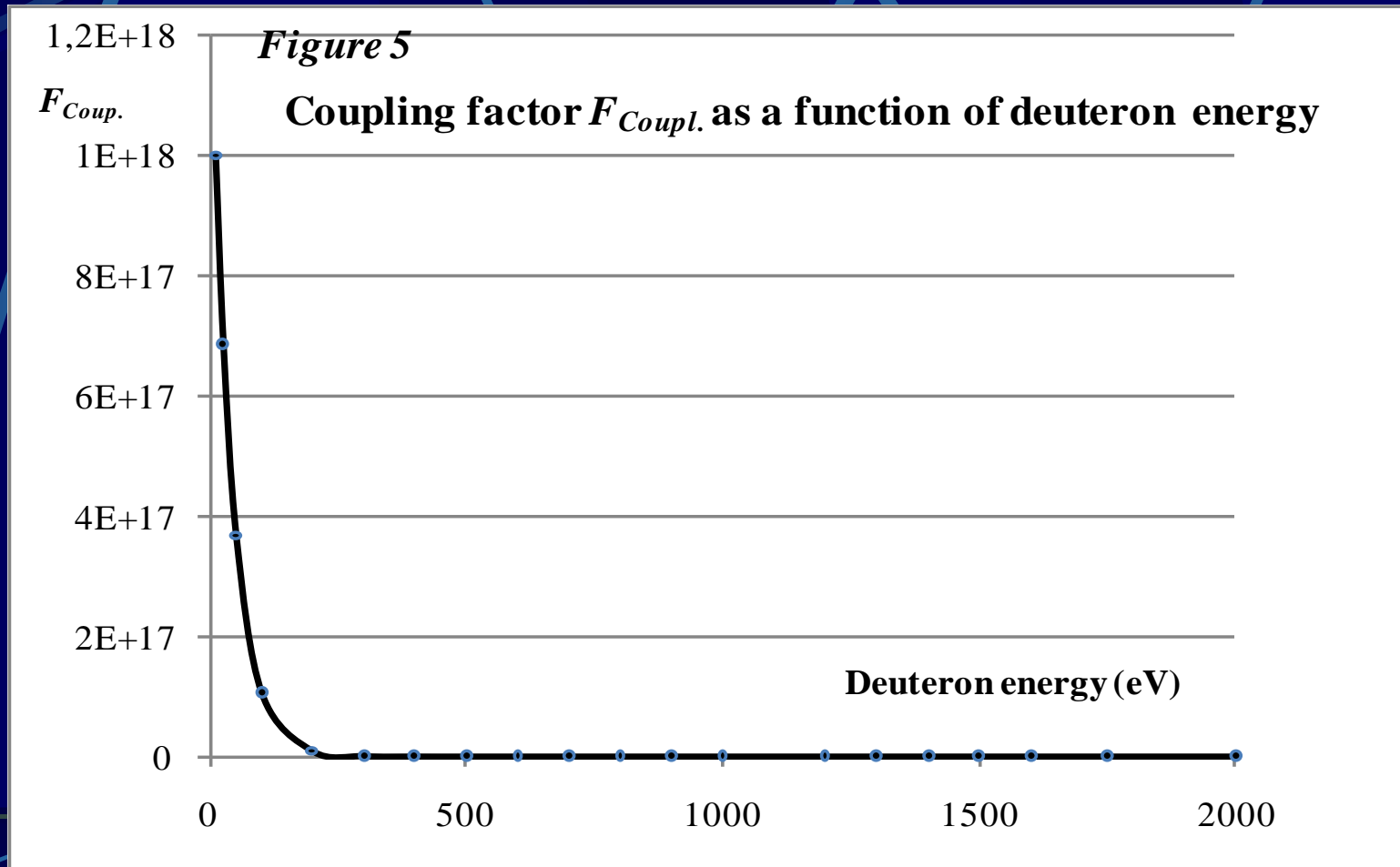
$$r_{ex} = F_c r_{cal}, \text{ with } F_c \cong 10^{18}$$

IV - Comparison experimental data versus calculated reaction rates

- Typical energies for Huccke experiments are 2000 to 10000 eV. For SPAWAR experiments they are round 2 eV
- The huge variation of F_c with E_d (energy of the deuteron) suggest a resonnant coupling between the impidging deuteron flux and the deuterons already present in the target (inducing fusion reactions between them).

$$F_{coupl.}(E_d) = F_{coupl.}(E_f) e^{-\Delta(F(E_f)) \frac{E_d - E_f}{E_f}} \quad \text{with } E_f \text{ fermi energy}$$

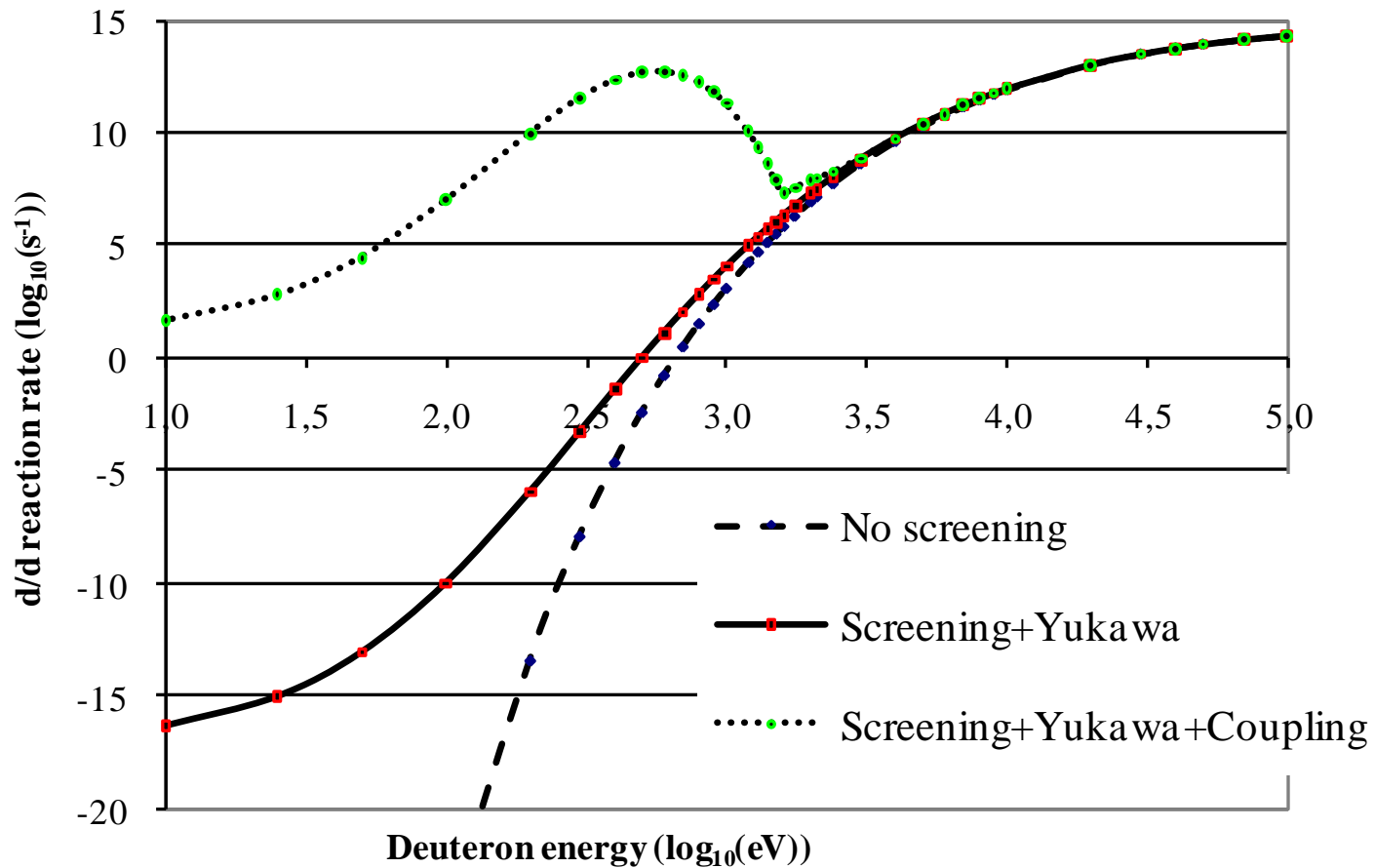
V - The coupling resonance and the CMAF window - The resonance curve



V - The coupling resonance and the CMAF window - d/d reaction rates

Figure 3

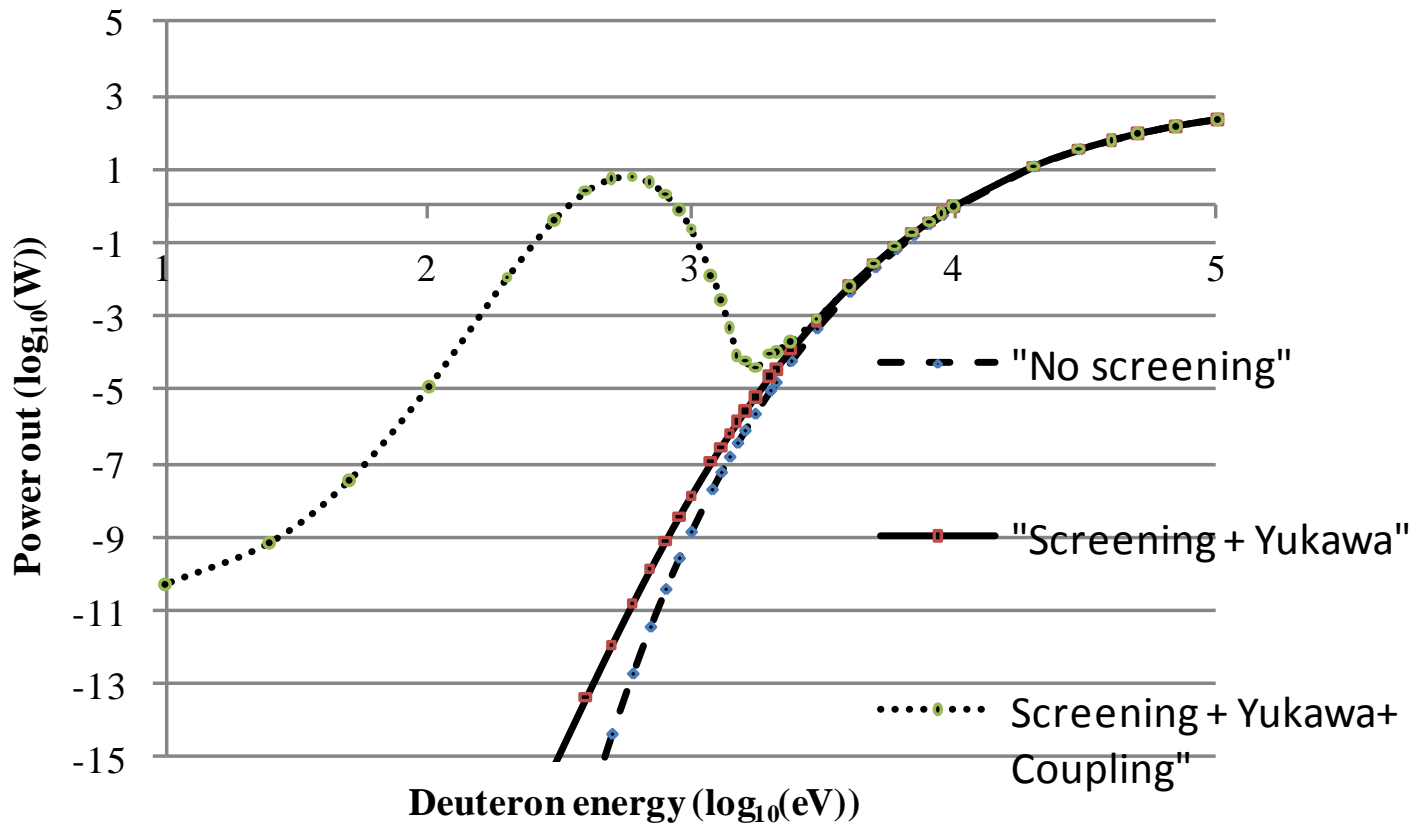
d/d reaction rate for 1 W in



V - The coupling resonance and the CMAF window - Energy out for 1 W in

Figure 4

Power out for 1 W in



V - The coupling resonance and the CMAF window - Energy out for 1 W in

- At low energy of the deuteron, the reaction rates are in line with SPAWAR results for a huge value of the coupling ($\#10^{18}$) The Yukawa potential has a negligible role. The energy production is very small (10^{-10} W). At a few keV, the Yukawa potential has an influence comparable to that of the electrons screening and with a coupling factor of 1.5 to 2, the results are in line with Huke measurements.
- At high energy of the deuteron, the reaction rates are in line with hot fusion.
- At energy of the deuteron between 500 and 1000 eV, sizeable energy production levels are expected. CMAF window (Condensed Matter Assisted Fusion)

V - The coupling resonance and the CMAF window - Production of deuterons beam in the CMAF window and target optimization

- Prototype for screening studies
- Ion guns for industrial applications
- Importance of F_c , depending upon physical and chemical characteristics of the target (optimization required).